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A new high-pressure phase of LiAlO₂

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Abstract

A new high-pressure phase of LiAlO₂ has been recovered through a shock recovery technique at pressures above 9 GPa. This new phase has been refined as a tetragonal structure with lattice parameters of a = 0.38866(8) nm and c = 0.83001(18) nm. Its calculated density is 3.51 g/cm^3 , about 34% denser than γ -LiAlO₂. The aluminum and lithium cations in this new phase are six-fold coordinated, as in α -LiAlO₂ and the structure of this new phase is similar to tetragonal LiFeO₂. This new high-pressure phase is stable at temperatures up to 773 K.

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1. Introduction

High-pressure technology is an important technique used to synthesize new dense materials. It enables us to produce a solid phase that is denser than the starting material. Although there are various high-pressure techniques, shock wave compression demonstrates a useful method to produce dense materials from their low-pressure phases. Examples include diamond from graphite [1], fullerites [2] and fullerene [3], super-hard cubic boron nitride from rhombohedral BN [4], cubic silicon nitride [5] and germanium nitride [6].

We have applied shock compression techniques to investigate the high-pressure phase transition in LiAlO₂. This type of lithium compound has potential applications in the energy industry as lithium battery cathodes [7]. LiAlO₂ has four known polymorphisms. The γ phase has a tetragonal symmetry with the space group $P4_12_12$ [8] and a density of 2.615 g/cm³. The orthorhombic β -phase has the same coordinated Li and Al cations as the γ -phase. The other two polymorphs can be synthesized from γ -phase under high pressures. One is the hexagonal α -phase with NaHF₂ structure synthesized at 3.5 GPa and 850°C [9]. Its density is 3.401 g/cm³. The other is a monoclinic phase synthesized at an intermediate pressure and temperature [10]. Sorting by coordination number, all the cations in the α -phase are coordinated by six oxygen anions and those in the γ phase by four oxygen anions. Lithium cations are sixfold coordinated and aluminum cations display both four- and six-fold coordination in the monoclinic phase. Experiments on LiAlO₂ at high pressures above 3.5 GPa have not been performed. It is expected that new dense high-pressure phases of LiAlO₂ may be stabilized at higher pressures.

In the present study, impact technology has been applied to compress γ -LiAlO₂ at pressures up to 48 GPa. We have succeeded in recovering a tetragonal, highpressure phase of LiAlO₂ with all cations six-fold coordinated by oxygen. This new phase has a distorted NaCl structure similar to the tetragonal LiFeO₂ phase. It is denser than the α -phase.

2. Experiments

 γ -LiAlO₂ powder with a purity higher than 99% was used in this study. The powder was used as received, and a mixture of γ -LiAlO₂ and copper powder (90 wt%) was used as the starting material. The starting materials were packed into steel containers. A propellant gun was used to accelerate a plastic projectile with a flyer plate (steel, aluminum or copper) to impact the container at the desired velocity. When the shock wave first enters the sample, a first shock condition of high pressure ($P_{\rm Hi}$)

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and high temperature $(T_{\rm Hi})$ is simultaneously generated. This shock condition reaches the equilibrium pressure $(P_{\rm Hf})$ and temperature $(T_{\rm Hf})$ after the shock wave has been reflected several times within the sample. This reflectance is due to the difference between the shock impedance of the sample and that of the container material. The recovered container was then mechanically cut open. Nitric acid solution was used to remove copper if a mixture was used as the starting material. Details of the recovery experiments can be found in Ref. [11]. Post-shock samples as well as the starting material were investigated by X-ray diffraction (XRD) using CuK α radiation, analytical electron microscopy (TEM), and electron energy loss spectroscopy (EELS).

3. Results and discussions

The shock conditions (P_{Hi} , T_{Hi} , P_{Hf} and T_{Hf}) are listed in Table 1 with other physical properties of the starting materials. The pressures, $P_{\rm Hi}$ and $P_{\rm Hf}$, and their corresponding temperatures for mixtures of LiAlO2 and copper were calculated with the shock impedance matching method [12] based on Hugoniot data [13] and thermodynamic constants of solid copper and porous copper. For shots on pure γ -LiAlO₂ powder, the shock temperatures and initial pressure were estimated based on thermal constants and Hugoniot data of sapphire. Sapphire was used for the estimate because of the lack of LiAlO2 data. This approximation may produce a 10% and 20% uncertainty in $P_{\rm Hi}$ and $T_{\rm Hi}$ for shots on pure LiAlO₂ and 5% and 10% uncertainty in pressure and temperature for the mixed powder, respectively.

XRD patterns for all the recovered samples of all the runs together with the starting γ -LiAlO₂ are illustrated in Fig. 1, where the runs 848, 852, 950, and 955 are all for a Cu/ γ -LiAlO₂ mixture (see Table 1). As can be seen in the figure, it appears that the intensities of some of the diffraction peaks decrease with increasing pressure and temperature while those of other peaks increase. The ones that increase correspond to the diffractions of a spinel type Al_{2.5}Li_{0.5}O₄ and its calculated lattice

Tab	le 1	l	
List	of	shock	conditions

constant is 0.7914(2) nm (from run 895), which is consistent with its reported value [14]. The peaks that decrease indicate the presence of a new phase, and they have been indexed as a tetragonal structure. The structure refinement by FullProf. program has been carried out with the Rietveld method [15]. The observed and calculated d values of the new phase in the run 950 are compared in Table 2, with lattice parameters of a =0.38866(8) and c = 0.83001(18) nm, four molecules of LiAlO₂, and its space group of $I4_1/amd$ (No. 141). The space group has been selected based on the known tetragonal LiFeO₂(18), and it is reasonable (Table 2). The density of this new phase is calculated to be 3.51 g/ cm^3 , about 34% denser than the y-phase. The atomic arrangements in real space of this new phase and spinel Al_{2.5}Li_{0.5}O₄ are illustrated in Fig. 2. A comparison is given in Fig. 3 by taking into account the coexistence of the spinel phase. Table 3 summarizes the refined structural data for the new phase. The calculated atomic occupancy suggests some defects and disordering in the new phase. In spinel Al_{2.5}Li_{0.5}O₄, Li atoms are coordinated by six oxygen atoms and Al by both six



Fig. 1. XRD results for γ -LiAlO₂, and recovered samples 848, 852, 895, 949, 950 and 955. Samples 895 and 949 were obtained with pure γ -LiAlO₂ powder while the others were obtained from a mixture of γ -LiAlO₂ and copper. Peaks for the new phase δ -LiAlO₂ are marked by arrows; the dot-marked bands in sample 895 indicate the overlapped diffractions of Li₂O and spinel.

Shot no.	Flyer material	Starting material	Impact velocity (km/s)	Sample density (g/cm ³)	Sample thickness (mm)	$P_{\rm Hi}$ (GPa)	$T_{\rm Hi}$ (K)	P _{Hf} (GPa)	$T_{\rm Hf}$ (K)
848	Steel	Mixture ^a	1.49	6.11	1.1	21	760	33	1100
852	Steel	Mixture ^a	1.07	6.50	1.8	15	600	22	780
895	Steel	LiAlO ₂	1.51	1.92	1.2	7.5	<350	35	<700
949	Copper	LiAlO ₂	2.00	2.24	0.5	10	<400	48	< 900
950	Al	Mixture ^a	0.77	6.43	1.8	6.2	420	9	500
955	Steel	Mixture ^a	1.99	6.48	1.9	30	1030	48	1570

^a Mixture of 10 wt% γ-LiAlO₂ and 90 wt% copper powder.

Table 2 X-ray diffraction data of a new phase δ -LiAlO₂

h k l		l	Observed d (nm)	Calculated d (nm)		
1	0	1	0.35123	0.35191		
0	0	4	0.20761	0.20765		
2	0	0	0.19387	0.19424		
2	1	1	0.16962	0.17006		
1	0	5	0.15245	0.15275		
2	0	4	0.14171	0.14186		
2	2	0	0.13717	0.13735		
2	1	5	0.12009	0.12007		
2	2	4	0.11442	0.11456		

This phase is indexed with a tetragonal symmetry with lattice constant a = 0.38866(8) nm and c = 0.83001(18) nm. The calculated density at room temperature is 3.51 g/cm^3 with Z = 4.



Fig. 2. Atom arrangements in spinel Al_{2.5}Li_{0.5}O₄ (a) and δ -LiAlO₂ (b). Oxygen, lithium, and aluminum atoms are illustrated by small (red), large (blue), and medium (yellow) spheres, respectively. It is obvious that all Li atoms in the spinel phase are six-fold coordinated, 40% of Al atoms are four-fold coordinated and 60% of Al atoms are six-fold coordinated (a) and both Li and Al are six-fold coordinated in δ -phase (b).

and four oxygen atoms, while all Li and Al are six-fold coordinated by oxygen in the new phase.

To characterize the chemical composition of the new phase, analytical electron transmission microscopy and EELS have been applied to sample 848. EELS spectra of lithium and aluminum in the selected area have been obtained. Based on EELS spectra of lithium and aluminum, the atomic number ratio of Li/Al was 1 ± 0.1 . According to the EELS data, the selected material should be LiAlO₂, because Al_{2.5}Li_{0.5}O₄ has a Li/Al ratio of 0.2. The K edge absorption of Li in γ -LiAlO₂, the recovered sample and LiF (as a reference) are illustrated in Fig. 4(a). The K absorption edges $(\sim 62 \,\text{eV})$ of the selected area and LiF are similar. This similarity shows that Li atoms in the new phase are sixfold coordinated. From Fig. 4(b), the similarity of the Labsorption edge of aluminum in the new phase to that in α -Al₂O₃ indicates that Al is six-coordinated, while the Al-L edge absorption of γ -LiAlO₂ indicates the presence of four-coordinated aluminum cations. The spectra for



Fig. 3. A comparison of XRD patterns. Plus (+) is the observed one, the overwritten curve is the calculated pattern, and the bottom curve indicates the difference between the two. Two series of vertical bars mark the Bragg positions of new phase δ -LiAlO₂ and spinel Al_{2.5}Li_{0.5}O₄, respectively.

Table 3								
Refined	structural	data for	the new	phase	δ -LiAlO ₂	with	conventi	onal
Rietveld	R-factor	of 12.9						

Atom	Wickoff	x	у	Ζ	В	Occ.
Li	4 <i>b</i>	0	0.75	0.625	2.19	0.814(13)
Al(2)	4b	0	0.75	0.625	1.5	0.186(13)
Al	4a	0	0.75	0.125	1.5	0.826(26)
Li(2)	4a	0	0.75	0.125	2.19	0.174(26)
0	8 <i>e</i>	0	0.75	0.34877(55)	1.00	1

Space group $I4_1/amd$ (No. 141), a = 0.38866(8) and c = 0.83001(18). Positions (x, y, and z), thermal parameter (B) of each atom, and site occupancy (occ.) are given together with their deviations.

spinel type Al₂O₃ differ from those of γ -LiAlO₂ and α -Al₂O₃. In conclusion, EELS spectra indicate that both Li and Al are four-fold coordinated in γ -LiAlO₂. This result differs from a nuclear magnetic resonance (NMR) study on γ -LiAlO₂, where ⁷Li was probed [16]. In that study lithium was found to be six-fold coordinated but according to the NMR shift of lithium oxide, shifts of ⁷Li are less than 10 ppm (Ref. [17]). This finding seems to be beyond the resolution of Ref. [16]. Therefore, the results of Ketty et al. on coordination number for y-LiAlO₂ may not be conclusive. The EELS data support the structure obtained by Marezio from single-crystal XRD data [8]. Furthermore our results clearly indicate that both Li and Al are coordinated by six oxygen atoms in the new phase. This result is consistent with the refined structure (Fig. 2b).

Based on the present XRD results and EELS spectra, a new high-pressure phase of LiAlO₂ was recovered. We call it δ -LiAlO₂. In order to obtain more information about the phase transition of γ -LiAlO₂ to δ -LiAlO₂ and



Fig. 4. (a) EELS spectra of Li-*K* in γ -LiAlO₂ (1), δ -LiAlO₂ (2) and LiF (3). (b) EELS spectra of Al- $L_{2,3}$ in γ -LiAlO₂ (1), δ -LiAlO₂ (2), α -Al₂O₃ (3), and γ -Al₂O₃ (4). XRD results for stability of δ -LiAlO₂ at high temperatures. Sample 950 with almost pure δ -LiAlO₂ was heated at 773, 843, 923 and 1073 K for 1 h, respectively, and the 923 K sample was further heated at 1173 K for 1 h.

phase stability of δ -LiAlO₂, shots done on pure γ -LiAlO₂ are depicted in Fig. 1 (runs 895 and 949). They are similar to the starting materials except for the addition of two broad diffraction lines, marked with dots in the figure. These broad bands correspond to the strongest diffractions of spinel Al_{2.5}Li_{0.5}O₄ and Li₂O. This can be explained as the following reaction:

$$5 \operatorname{LiAlO}_2 = 2 \operatorname{Al}_{2.5} \operatorname{Li}_{0.5} \operatorname{O}_4 + 2 \operatorname{Li}_2 \operatorname{O}_2.$$

Spinel Al_{2.5}Li_{0.5}O₄ can result from the decomposition of LiAlO₂, instead of transition to δ -LiAlO₂, under shock-induced high-temperature and high-pressure conditions, although the product Li₂O was removed during the treatment by nitric acid of samples 848, 852, 950 and 955. From Table 1, shots 895 and 949 have a very low initial shock pressures and low shock temperatures because pure γ -LiAlO₂ was used as a starting material, and these samples do not indicate the transition of γ - to δ -LiAlO₂. Copper powder can act as an efficient agent of both heating samples and enhancing pressures. The recovered powder 949 was observed to be only slightly sintered while sample 895 was well sintered, indicating that a higher temperature was generated in shot 895 than in shot 949. At enough high temperatures, LiAlO₂ was decomposed to spinel Al_{2.5}Li_{0.5}O₄. So spinel Al_{2.5}Li_{0.5}O₄ is a byproduct from high temperature decomposition of γ -LiAlO₂, while δ -LiAlO₂ will be a product of high-pressure solid–solid phase transition. Sample 955, which experienced the highest peak pressure of 48 GPa, is found to contain a certain amount of δ -phase (Fig. 1). This indicates that δ -LiAlO₂ is still stable at shock pressures up to 48 GPa.

It is worthy to address the structural similarity between LiAlO₂ and LiFeO₂. LiFeO₂ has three polymorphisms (with rhombohedral, cubic, and tetragonal symmetry, respectively) of NaCl or distorted NaCl structure, in which both Li and Fe are six-fold coordinated by oxygen. LiAlO₂ has five polymorphisms, a monoclinic phase, rhombohedral *a*-phase, orthorhombic β -phase, tetragonal γ - and δ -phases. In general, the radius of Al^{3+} ion in LiAlO₂ is close to that of Fe³⁺ in LiFeO₂. α -LiAlO₂ is isostructural to the rhombohedral LiFeO2. According to solid-state chemistry [18], isostructurality can be expected between phases of LiAlO₂ and LiFeO₂, but the previously known tetragonal γ -LiAlO₂ does have a different structure to that of tetragonal LiFeO2 as seen in different coordination environments of the cations. However, our obtained δ -LiAlO₂ appears to be isostructural to tetragonal LiFeO₂. The structural similarity between α -LiAlO₂ and rhombohedral LiFeO₂, and that between δ -LiAlO₂ and tetragonal LiFeO₂ indicates that high-pressure phases of LiFeO₂ can shed some light on the possible high-pressure structure of LiAlO₂, beyond the δ -phase. These possible candidates may have a CsCl type or distorted CsCl structure.

We have also investigated the temperature stability of δ -LiAlO₂. Several samples taken from sample 950, which consists of nearly pure δ -LiAlO₂, were heated in the atmosphere at 773, 843, 923, 1073 and 1173 K for 1 h, respectively. The XRD of each heated sample is shown in Fig. 5. δ -LiAlO₂ is stable up to 773 K. When heated to 843 K, a certain amount of a-phase was detected. With increasing temperatures up to 923 K, most of the δ -phase turned into α -phase. A little γ -phase appeared when the sample was heated to 1073 K. The sample heated at 1173 K produced the γ -phase. These heating experiments indicate a series of reversed transformations to the low-pressure phase with increasing temperature. The phase transitions of LiAlO₂ from δ - to α -phase and from α - to γ -phase reinforce that the composition of δ -phase has the same chemical composition as $LiAlO_2$. The phase-transition temperature from



Fig. 5. XRD results for stability of δ -LiAlO₂ at high temperatures. Sample 950 with almost pure δ -LiAlO₂ was heated at 773, 843, 923, and 1073 K for 1 h, respectively, and the 923 K sample was further heated at 1173 K for 1 h. XRD measurements were performed at ambient conditions.

 α - to γ -LiAlO₂ is about 1073 K, which is consistent with reported results [19].

It is interesting to examine the pressure ranges when α - and δ -phase are stable. According to the heating results mentioned above, it appears that δ -LiAlO₂ is a high-pressure stable phase compared to α -LiAlO₂. This is because an increase in temperature always tends to decrease the density of a solid while a rise in pressure increases it. However, when LiFeO₂ is considered, its temperature stability is different. LiFeO₂ with rhombohedral symmetry is a low-temperature phase while LiAlO₂ with tetragonal symmetry is the low-temperature case. Because the pressure ranges from 3 to 9 GPa is out of our experimental capability, we cannot get direct information on the possibility of synthesizing α -LiAlO₂ by shock compression. It may be useful to increase static high pressures applied to LiAlO₂ up to 10 GPa to investigate the possibility of synthesizing this new highpressure phase. This technique is currently out of our research range.

It should be interesting to extend the same shock recovery experiments of LiAlO₂ to its analogue compounds, such as LiBO₂, NaAlO₂ and LiGaO₂. Several shock compressions on pure powders or powders mixed with copper have been conducted for each oxide mentioned above in the pressure ranges of 10-60 GPa and temperature ranges of 700-2500 K. No new phase was successfully detected by the XRD investigation on the recoveries. Some attention should be paid to the following observations: (1) phase-structure similarity between LiAlO₂ and LiGaO₂ at pressures up to 3.5 GPa [10] may not be extrapolated to the higher-pressure

range, and (2) a long-term diffusive process, not realized by shock compression, may play an important role in these compounds.

4. Conclusions

A new high-pressure phase of LiAlO₂ (δ -phase), in which both lithium and aluminum are six-fold coordinated, has been synthesized from shock compressed γ -LiAlO₂. It is denser than α -phase. This new δ -phase is stable when shocked to 48 GPa at 1600 K. It is also stable when heated in the atmosphere at 773 K. δ -LiAlO₂ is isostructural to tetragonal LiFeO₂, while an isostructure to rhombohedral LiFeO₂ is α -LiAlO₂.

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